Lecture 2: Linear Systems - Part 1 Winter 2016

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Linear time-invariant systems, Linear ODEs, and the power of decomposability



Joseph-Louis LaGrange 1736 - 1813



Pierre-Simon LaPlace 1749 - 1827



Joseph Fourier 1768 - 1830 So, today we continue to stay with relatively small-scale linear systems...

	n = 1	n = 2 or 3	n >> 1	continuum
Linear	exponential growth and decay single step conformational change fluorescence emission pseudo first order kinetics	second order reaction kinetics linear harmonic oscillators simple feedback control sequences of conformational change	electrical circuits molecular dynamics systems of coupled harmonic oscillators equilibrium thermodynamics diffraction, Fourier transforms	Diffusion Wave propagation quantum mechanics viscoelastic systems
Nonlinear	fixed points bifurcations, multi stability irreversible hysteresis overdamped oscillators	anharmomic oscillators relaxation oscillations predator-prey models van der Pol systems Chaotic systems	systems of non- linear oscillators non-equilibrium thermodynamics protein structure/ function neural networks the cell ecosystems	Nonlinear wave propagation Reaction-diffusion in dissipative systems Turbulent/chaotic flows

So, for today, linear systems analysis...

The goals will be three-fold:

(1) Understand the origins if simplicity in linear, time-invariant systems (the general regime of most modern engineering)

(2) Understand the principle of decomposability of linear systems...using a simple model of a second order process

(3) Learn a new way to solve differential equations that makes the concepts of linearly and decomposability more intuitive...

First let's consider a change of perspective for today about how we regard basic reactions. We will change from the usual state-contric to a process - rentric vew :

$$\begin{array}{c} A \cdot Folding \\ A \cdot Foldi$$

$$\frac{\mathbb{B}}{\operatorname{hv}} \xrightarrow{\operatorname{cond}} \mathbb{A} \xrightarrow{\operatorname{dv}} \mathbb{M} = \mathbb{P} \quad S(\mathcal{H}) \longrightarrow [\mathcal{H}(\mathcal{H})] \longrightarrow \mathbb{M}(\mathcal{H})$$

$$\stackrel{\operatorname{tor}}{\operatorname{tor}} \mathcal{R}_{0} \xrightarrow{\operatorname{tor}} \mathbb{R}_{0}$$

$$\stackrel{\operatorname{hv}}{\operatorname{hv}} \xrightarrow{\operatorname{so}} \mathcal{R} \rightarrow \mathbb{I}_{1} \rightarrow \mathbb{I}_{2} + \mathbb{I}_{3} \rightarrow \mathbb{M} = \mathbb{P} \quad S(\mathcal{H}) \rightarrow [\mathcal{H}(\mathcal{H})] \rightarrow [\mathcal{H}_{2}(\mathcal{H})] \rightarrow \cdots$$

$$\stackrel{\operatorname{hv}}{\operatorname{hv}} \xrightarrow{\operatorname{so}} \mathcal{R} \rightarrow \mathbb{I}_{1} \rightarrow \mathbb{I}_{2} + \mathbb{I}_{3} \rightarrow \mathbb{M} = \mathbb{P} \quad S(\mathcal{H}) \rightarrow [\mathcal{H}(\mathcal{H})] \rightarrow [\mathcal{H}_{2}(\mathcal{H})] \rightarrow \cdots$$

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First let's consider a change of perspective for today about how we regard busic reactions. We will change from the usual state-contric to a process - centric vew :



...this is a basic **feedback control circuit**; a function of the output is fed back to the input to clamp output levels.



...for example, a system for constant-density growth of micro-organisms (a so-called "**turbidostat**")

How can we understand/ predict the behavior of those systems ? well, it the systems are linear and time invariant, there is a powerful theory Sur. () definitions (3) properties of livearity and time invanance (3) convolution theorem (9) The haplace transform method [ for simplicity and intuition] (5) Three examples first order second order -1-13simple feedback system.

...so the theory of LTI systems.

#### Part 1: Linear Time-Invariant Systems

A theme will be to understand the "simplicity" of linear, time-invariant systems. What does "linearity" and "time-invariance" mean exactly?

$$\begin{split} & \text{Intervity implies a principle ralled superposition : If  $y_1(d)$  to the output of a system to imput  $y_1(d)$  and  $y_2(d)$  is the response to  $y_3(d)$ . Then:  

$$\begin{split} & \textcircled{O} \quad x_1(d) + y_2(d) \longrightarrow \quad y_1(d) + y_2(d) \qquad (additionly] \\ & \textcircled{O} \quad a_1 \cdot x_1(d) \longrightarrow \quad a_1 \cdot y_1(d) \qquad (additionly] \\ & \textcircled{O} \quad a_1 \cdot x_1(d) \longrightarrow \quad a_2 \cdot y_1(d) \qquad (scaling or homogonally) \\ & \overbrace{O} \quad a_1 \cdot x_1(d) \longrightarrow \quad a_2 \cdot y_1(d) \qquad (scaling or homogonally) \\ & \overbrace{O} \quad a_1 \cdot x_1(d) \longrightarrow \quad a_2 \cdot y_1(d) \qquad (scaling or homogonally) \\ & \overbrace{O} \quad a_1 \cdot x_1(d) \longrightarrow \quad a_2 \cdot y_1(d) \qquad (scaling or homogonally) \\ & \overbrace{O} \quad a_1 \cdot x_1(d) \longrightarrow \quad a_2 \cdot y_1(d) \qquad (scaling or homogonally) \\ & \overbrace{O} \quad a_1 \cdot x_1(d) \longrightarrow \quad a_2 \cdot y_1(d) \qquad (scaling or homogonally) \\ & \overbrace{O} \quad a_1 \cdot x_1(d) \longrightarrow \quad a_2 \cdot y_1(d) + a_2 \cdot y_2(d) + \cdots \\ & \overbrace{V} \quad (d) : \quad \underbrace{\sum_{k} a_k \cdot y_k(d)}_{k \in \mathbb{N}} = a_1 \cdot y_1(d) + a_2 \cdot y_2(d) + \cdots \\ & \overbrace{V} \quad (d) : \quad \underbrace{\sum_{k} a_k \cdot y_k(d)}_{k \in \mathbb{N}} = a_1 \cdot y_1(d) + a_2 \cdot y_2(d) + \cdots \\ & \overbrace{V} \quad (d) : \quad \underbrace{\sum_{k} a_k \cdot y_k(d)}_{k \in \mathbb{N}} = a_1 \cdot y_1(d) + a_2 \cdot y_2(d) + \cdots \\ & \overbrace{V} \quad (d) : \quad \underbrace{\sum_{k} a_k \cdot y_k(d)}_{k \in \mathbb{N}} = a_1 \cdot y_1(d) + a_2 \cdot y_2(d) + \cdots \\ & \overbrace{V} \quad (d) : \quad \underbrace{\sum_{k} a_k \cdot y_k(d)}_{k \in \mathbb{N}} = a_1 \cdot y_1(d) + a_2 \cdot y_2(d) + \cdots \\ & \overbrace{V} \quad (d) : \quad \underbrace{\sum_{k} a_k \cdot y_k(d)}_{k \in \mathbb{N}} = a_1 \cdot y_1(d) + a_2 \cdot y_2(d) + \cdots \\ & \overbrace{V} \quad (d) : \quad \underbrace{\sum_{k} a_k \cdot y_k(d)}_{k \in \mathbb{N}} = a_1 \cdot y_1(d) + a_2 \cdot y_2(d) + \cdots \\ & \overbrace{V} \quad (d) : \quad \underbrace{\sum_{k} a_k \cdot y_k(d)}_{k \in \mathbb{N}} = a_1 \cdot y_1(d) + a_2 \cdot y_2(d) + \cdots \\ & \overbrace{V} \quad (d) : \quad \underbrace{\sum_{k} a_k \cdot y_k(d)}_{k \in \mathbb{N}} = a_1 \cdot y_1(d) + a_2 \cdot y_2(d) + \cdots \\ & \overbrace{V} \quad (d) : \quad \underbrace{\sum_{k} a_k \cdot y_k(d)}_{k \in \mathbb{N}} = a_1 \cdot y_1(d) + a_2 \cdot y_2(d) + \cdots \\ & \overbrace{V} \quad (d) : \quad \underbrace{\sum_{k} a_k \cdot y_k(d)}_{k \in \mathbb{N}} = a_1 \cdot y_1(d) + a_2 \cdot y_2(d) + \cdots \\ & \overbrace{V} \quad (d) : \quad \underbrace{\sum_{k} a_k \cdot y_k(d)}_{k \in \mathbb{N}} = a_1 \cdot y_1(d) + a_2 \cdot y_2(d) + \cdots \\ & \overbrace{V} \quad (d) : \quad (d$$$$

These two properties underlie the concept of the convolution & integral ...

Stept: The impulse function. The first concept we need is that of  
representing any arbitrary input function as a series  
of scaled, time shifted elementary functions called  
impulses...  
Su discrete time...  

$$\delta[n] = \begin{cases} 0 & n \neq 0 \\ 2 & n = 0 \end{cases}$$
  
 $\int (n-2)$ 

Step1: The impulse function. The first concept we used is that g  
representing any arbitrary input function as a sence  
of scaled, time shifted elementary fluctums called  
impulses...  
In continues time...  
$$\frac{1}{100} = \frac{200}{100} = \frac{100}{100} = \frac{200}{100} = \frac{100}{100} = \frac{100}{100$$









This defines the so-called "impulse response" of our system....the output due to an impulse stimulus.



This defines the so-called "impulse response" of our system....the output due to an impulse stimulus.

What then will be there system response to some arbitrary function?

Step 11 wow we give new system an input function ... you X(m) X(ch) 0.5 0 h xio)h(o) × (0) \$(0) × 0.5 0.5 -1 0 -2 ۰. 6 x(1)k(n-1) X(1) \$(n-1) 2 5 4 012 10 1- 5ų. 25 25 yen) 0.5 . . . . . y(n)= x(0)h(0) + x(1) h(n-1) 50 ....

Morr generally ...  

$$y(n) = \sum_{k=-\infty}^{\infty} x(k) h(n-k) \qquad Have is the convolution sum or
supreposetion sum$$
For LTI systems , this gives the output (y(n)) to any arbitrary reput (xo  
given knowledge of its impulse response (h(n)).
  
Thus is a major result of linear systems analysis ... the response of an  
LTI system is completely defined by its impulse response.

More generally ...  

$$y(n) = \sum_{k=-\infty}^{\infty} x(k) h(n-kk) \qquad \text{there is the convolution summ}$$

$$for the systems for any arbitrary upst (xind)
given knowledge of its impulse response (h(n)).
There is a major result of linear systems analysis ... the response of an
$$b Ti system is completely defined by its impulse response.$$
In continuous time ...
$$x(t) \rightarrow h(t) \rightarrow y(t)$$$$

h(+)

x(+)

The definition of the convolution operator...and the power of LTI systems....you can completely characterize them by their response to the simplest input...the impulse response! Now...a simple graphical way of understanding convolution...the sliding of two functions across each other...



Now...a simple graphical way of understanding convolution...the sliding of two functions across each other...





A defferential expension is an equation with derivatives in it.  
In the general case:  
$$g\left[f(4), \frac{44}{44}, \frac{14}{44^2}, \dots, \frac{d^2f}{de^2}; t\right] = h(t)$$

to refresh...as nicely explained in the mathematics course yesterday

**K. A. Reynolds**, "Mathematical Foundations", lecture 1

But ... the simplest case is a linear, 1<sup>st</sup> order, homogonrous one:  

$$dy = f(t,y)$$
 (Homogenen, 1<sup>st</sup> order, homogonrous one:  
 $dy = f(t,y)$  (Interver, 2<sup>st</sup> order, homogonrous one:  
 $dy = f(t,y)$  (

**K. A. Reynolds**, "Mathematical Foundations", lecture 1



...so this is a second order, linear, homogeneous equation...the equation of motion for **a** harmonic oscillator

$$\frac{d^2x}{dt^2} + \alpha \frac{dx}{dt} + \beta x = 0 \qquad [110000 gaussus, 100000, 200 der]$$

Note that our second order system just got reduced to **two first order** differential systems!! The solution to the linear, first-order, homogeneous differential equation...

This is the availytic solution.

again...as explained in the mathematics course yesterday

For example, both binding and dissociation reactions for bimolecular interaction at the pseudo-first order limit is well described by a first-order process...





What about a second order system...the case of multistep conformational change...



#### <u>Rhodopsin</u>

...a **system** of differential equations. What is the solution for **A(t)** and **B(t)**?

What about a **second order system**...the case of multistep conformational change...

i) 
$$\frac{dA}{dt} = -k_{0}A$$
  
i)  $\frac{dB}{dt} = k_{0}A - k_{0}B$   
BG

...the solutions:

$$A(t) = A_0 e^{-kt}$$

$$B(t) = \frac{A_0 k_1}{k_2 - k_1} \left[ e^{-k_1 t} - e^{-k_2 t} \right]$$

...as demonstrated in the mathematics course yesterday
The solution to the linear, first-order, inhomogeneous differential equation...

.

$$\frac{dB}{dt} = -k_s B + k_s A, \quad where \quad A = A_s e^{-k_s t}$$

$$B(s) = 0$$

**K. A. Reynolds**, "Mathematical Foundations", lecture 1

The solution to the linear, first-order, inhomogeneous differential equation...

$$\frac{dB}{dt} = -k_2B + k_1A, \quad \text{where} \quad A = A_0e^{-k_1t}$$

$$B(0) = 0$$

, bunch of algebra, using initial conditions...

.

$$B(+) = \frac{A_0 k_1}{k_2 - k_1} \left[ e^{-k_1 t} - e^{-k_2 t} \right]$$

Laplace transforms...

An approach to solving such equations that leads to **some important intuition** about linear systems...

Laplace transforms...

Given any continuously differentiable function 
$$f(t)$$
, we define the  
Laplace transform of  $f(t)$ :  
 $d\left[f(t)\right] = F(s) = \int_{0}^{\infty} f(t) e^{-st} dt$   
Sec. this is a transformation of  $f(t)$  so that:  
 $f(t) \stackrel{\checkmark}{\longleftarrow} F(s) = \int_{0}^{\infty} f(t) = \int_{0}^{\infty} f(t) = \int_{0}^{\infty} f(t) \int_{0}^{\infty} f$ 

Why should we do this transformation?

Laplace transforms...

sme lace transforms: met  $( f e ) = e^{-kt} :$   $\int \{fe_{i}\} =$ Then ....

Some useful haplace transforms:  
(D) 
$$f(t) = e^{-kt}$$
; Then...  
 $J \{f(t)\} = F(s) = \int_{0}^{\infty} e^{-kt} e^{-st} dt$   
 $= \int_{0}^{\infty} e^{-(k+s)t} dt$   
 $= -\frac{1}{k+s} e^{-(k+s)t} \int_{0}^{\infty}$   
 $= \frac{1}{s+k}$ 

Some useful hapling transforms:  
(1) 
$$f(t) = e^{-kt}$$
: Then...  
 $J \{f(t)\} = F(s) = \int_{0}^{\infty} e^{-kt} e^{-st} dt$   
 $= \int_{0}^{\infty} e^{-(k+s)t} dt$   
 $= -\frac{1}{k+s} e^{-(k+s)t} \int_{0}^{\infty}$   
 $= \frac{1}{s+k}$ 

Some 
$$\frac{f(t)}{e^{-kt}} = \frac{f(s)}{s+ke}$$

(a) f(t) = t; Then...  $F(s) = \int_{0}^{\infty} t e^{-st} dt$ 

(2) f(4) = t ; Then ... F(s)= footest dt To solve this, we use the old integrate by parts rule: Judy = uv = Judu I've own case, set: u=t dv= estdt  $\int te^{st} dt = -\frac{t}{s}e^{-st} \int_{0}^{\infty} - \int_{0}^{\infty} -\frac{1}{s}e^{-st} dt$  $O = \left(\frac{1}{s^2} e^{-st}\right) \Big|_{0}^{\infty}$ 

(a) 
$$f(t) = t$$
; Then...  

$$F(s) = \int_{0}^{\infty} t e^{-st} dt$$
So ....  

$$\frac{f(t)}{t} = \frac{F(s)}{s^{2}}$$

(3) Now ... 
$$f(4) = \frac{d4}{dt}$$
  

$$\int \left\{ \frac{dt}{dt} \frac{d4}{dt} \right\} = \int_{0}^{\infty} \frac{d4}{dt} e^{-st} dt$$

$$\lim_{t \to \infty} \frac{d4}{dt} \frac{d4}{dt} = \int_{0}^{\infty} \frac{d4}{dt} e^{-st} dt$$

$$\lim_{t \to \infty} \frac{d4}{dt} \frac{d4}{dt} \frac{d4}{dt}$$

$$\int_{0}^{\infty} \frac{d4}{dt} e^{-st} dt = -f(t) e^{-st} \Big|_{0}^{\infty} - \int_{0}^{\infty} f(t) (-s e^{-st}) dt$$

$$= -f(0) + s \int_{0}^{\infty} f(t) e^{-st} dt$$

$$= -f(0) + s F(s)$$
So ....  $\frac{f(4)}{dt} = \frac{f(4)}{s} \frac{f(s)}{s}$ 

and  $\dots$  where proof (1) f(t) F(s) $a_1f_1(t) + a_2f_2(t)$   $a_1F_1(s) + a_2F_2(s)$ 

Oten-using all this, we go back to our deferential equation:  

$$\frac{dA}{JE} = -kA, \quad h(b) = A_0$$

$$J \left\{ \frac{dA}{JE} \right\} = J \left\{ -kA \right\}$$

$$SA(s) = A(0) = -kA(s)$$

$$SA(s) = kA(s) = A_0$$

$$A(s) = \frac{A_0}{s+k}$$
The heavilow solution.

To find our solution, we just take the month haplace transform:  
A(s)= 
$$\frac{4\pi}{3\pi k}$$
  
 $\chi^{-1}(A(s)) = \chi^{-1}[\frac{A_0}{3\pi k}]$ 

To find our solution, we just take the momenter haplace transform:  
A(s): 
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remember	

So	ተው)	F(5)	
	e <sup>-kt</sup>	stle	

To find our solution, we just take the moment haplace transform:  

$$A(s) = \frac{4s}{s+k}$$

$$J^{-1}(A(s)) = J^{-1} \left[ \frac{A_0}{s+k} \right]$$

$$A(t) = A_0 e^{-kt}$$
Note that we didn't integrate anything ! we just did algebra.

In Laplace transform space (F(s)), differentiation and integration become just a matter of doing algebra and finding the inverse transform...

What about the more complicated inhomogeneous first order equation?

$$\frac{dg}{dt} = -k_2 B + k_1 A, \quad where \quad A = A_0 e^{-k_1 t}$$

$$B(o) = 0$$

The solution to the **linear**, **first-order**, **inhomogeneous** differential equation... the old way...

$$\frac{dB}{dt} = -k_2B + k_1A, \quad \text{where} \quad A = A_0 e^{-k_1^2}$$

$$B(0) = 0$$

bunch of algebra, using initial conditions...

.

$$B(+) = \frac{A_0 k_1}{k_2 - k_1} \left[ e^{-k_1 t} - e^{-k_2 t} \right]$$

 $\frac{de}{dt} = -k_{x}B + k_{x}A$   $\frac{dB}{dt} + k_{x}B = k_{x}A_{0}e^{-k_{x}t}$   $Z\left[\frac{dB}{dt}\right] + k_{x}A\left\{8\right\} = k_{x}A_{0}Z\left\{e^{-k_{x}t}\right\}$ 

$$\frac{dg}{dt} = -k_{x}G + k_{x}A$$

$$\frac{dg}{dt} + k_{x}G = k_{x}A_{0}e^{-k_{x}t}$$

$$\frac{dg}{dt} + k_{x}A_{0}g = k_{x}A_{0}e^{-k_{x}t}$$

$$\frac{f(dg}{dt}) + k_{x}A_{0}g = k_{x}A_{0}\int e^{-k_{x}t}$$

$$SG(x) = B(x) + k_{x}B(x) = k_{x}A_{0}\int \frac{1}{3+k_{x}}$$

$$B(x) = 5 + k_{x} = k_{x}A_{0}\int \frac{1}{3+k_{x}}$$

$$B(x) = \frac{k_{x}A_{0}}{(3+k_{x})(x+k_{x})}$$
The Tangent solution

$$\frac{dg}{dt} = -\frac{k_{x}G}{k_{x}G} + \frac{k_{y}A}{k_{y}G}$$

$$\frac{dg}{dt} = -\frac{k_{x}G}{k_{x}G} + \frac{k_{y}A}{k_{y}G} = \frac{k_{y}A}{k_{y}G}$$

$$\frac{dg}{dt} + \frac{k_{x}}{k_{y}} \left\{ \frac{dg}{dt} \right\} + \frac{k_{x}}{k_{y}} \left\{ \frac{dg}{dt} \right\} = \frac{k_{y}A}{k_{y}} \left\{ \frac{dg}{dt} - \frac{k_{z}A}{k_{y}} \right\}$$

$$SG(4) = B(6) + \frac{k_{z}B(5)}{k_{z}} = \frac{k_{y}A}{k_{z}} \left[ \frac{1}{s+k_{y}} \right]$$

$$B(6) \left[ \frac{s}{s+k_{z}} \right] = \frac{k_{z}}{k_{z}} A_{0} \left[ \frac{1}{s+k_{y}} \right]$$

$$B(6) \left[ \frac{s}{s+k_{z}} \right] = \frac{k_{z}}{k_{z}} A_{0} \left[ \frac{1}{s+k_{y}} \right]$$

$$The Transform solution
$$\int_{0}^{1} \left\{ \frac{g(2)}{k_{z}} \right\} = \frac{k_{z}A}{k_{z}} \int_{0}^{1} \left\{ \frac{1}{(s+k_{z})(s+k_{z})} \right\}$$

$$\# she holds$$$$

	A Short Table of Laplace Transforms					
	y = f(t), t > 0 [ $y = f(t) = 0, t < 0$ ]	$Y = L(y) = F(p) = \int_0^\infty e^{-pt} f(t) dt$				
<i>L</i> 1	1	1	Rc p > 0			
₹ L2	r-4	$\frac{1}{p+a}$	$\operatorname{Re}(p+a) > 0$			
L3	sin at	$\frac{a}{p^2 + a^2}$	Re p > [1m a]			
14	COS al	$\frac{p}{p^2 + a^2}$	Re p≥  lm a			
LŚ	$t^{k}, k > -1$	$\frac{k!}{p^{k+1}}  \text{or}  \frac{\Gamma(k+1)}{p^{k+1}}$	Re p > 0			
L6	$t^k e^{-at},  k > -1$	$\frac{k!}{(p+a)^{k+1}}$ or $\frac{\Gamma(k+1)}{(p+a)^{k+1}}$	$\operatorname{Re}(p+a) > 0$			
> 17	$\frac{e^{-at}-e^{-bt}}{b-a}$	$\frac{1}{(p+a)(p+b)}$	$\operatorname{Re}(p+a) > 0$			
LS	$\frac{ae^{-ae}-be^{-be}}{a-b}$	$\frac{p}{(p+a)(p+b)}$	and Re $(p + \delta) > 0$			

Lookup inverse transforms...

$$\frac{dg}{dt} = -k_{z}g + k_{z}A$$

$$\frac{dg}{dt} = -k_{z}g + k_{z}A$$

$$\frac{dg}{dt} = +k_{z}B = k_{z}A_{z}e^{-k_{z}t}$$

$$\frac{dg}{dt} = +k_{z}A\{g\} = k_{z}A_{z}e^{-k_{z}t}\}$$

$$\frac{dg}{dt} = +k_{z}A\{g\} = k_{z}A_{z}e^{-k_{z}t}\}$$

$$\frac{dg}{dt} = -g(e) + k_{z}g(s) = k_{z}A_{z}\left[\frac{1}{s+k_{z}}\right]$$

$$\frac{dg}{dt} = \frac{A_{z}K_{z}}{k_{z}-k_{z}}\left[\frac{1}{s+k_{z}}\right]$$

$$\frac{dg}{dt} = -k_{z}g + k_{z}A$$

$$\frac{dg}{dt} = -k_{z}g + k_{z}A$$

$$\frac{dg}{dt} = +k_{z}B = k_{z}A_{z}e^{-k_{z}t}$$

$$\frac{dg}{dt} = +k_{z}A\{B\} = k_{z}A_{z}e^{-k_{z}t}\}$$

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$$\frac{dg}{dt} = -g(e) + k_{z}g(e) = k_{z}A_{z}\left[\frac{1}{s+k_{z}}\right]$$

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$$\frac{dg}{dt} = -g(e) + k_{z}g(e) = k_{z}A_{z}A_{z}\left[\frac{1}{s+k_{z}}\right]$$

$$\frac{dg}{dt} = k_{z}A_{z}A_{z}\left[\frac{1}{s+k_{z}}\right]$$

Much simpler approach...

Part 3: Back to decomposability of linear systems...

But...also an approach to "see" the reduction of a high-order system to a combination of first-order systems...

In keeping with our process-centric view, this of a first order system as a "process" that converts an input into an output.

The process has a characteristic function...the so-called "transfer function". It is fundamentally defined by the **response to an impulse input**.

Thus, outputs are <u>predictable</u> from just convolutions of the impulse response with the input function...

Lets choole ner unifierstranding....  
Thus, a second order system is like linking two first order systems  
type Theor, where the output of the first is the imput to the  
second:  
Typet 
$$(4, (4))$$
  $(4, (4))$   $(4, (4))$   $(4, (4))$   $(3, (4))$   $(4, (4)) = e^{-k_1 t}$   
 $(1, (4))$   $(4, (4))$   $(4, (4))$   $(4, (4))$   $(3, (4))$   $(1, (4)) = e^{-k_1 t}$   
 $(1, (4))$   $(1, ($ 

Let's use this to study our two cases...a first order system (single exponential) and our second order system (a double exponential)...



So, the first order system can be thought of as a system driven by an impulse stimulus...



Constrained and on processo of producting 
$$B(t)$$
 is due to the seried of the function  
of first order processo is a seried of the function of the function is the seried of the function of t

So, a second order system is a series of two first order systems....a basic property of linearity!

This reduction of a high-order system to a combination of first-order systems is a fundamental property of linear systems...

Is this general for any order equation? Yes...  
Soyue have:  

$$\frac{d!^{n}u}{dx^{n}} = f(t_{j}u) \qquad (n^{t_{j}} \text{ order })$$
we can decompose this into a system of a first order equa:  

$$\frac{dq_{inten}}{dt} = \frac{d}{dt} = \frac{dq_{i}}{dt} = \frac{dq_{i}}{$$

So...linear time-invariant systems are "simple" (not complex) for two reasons:

(1) They have the property that the impulse response fully characterizes their behavior. All responses to more complex inputs are just a convolution of the impulse response (the transfer function) with the input function.

Input 
$$\rightarrow$$
  $H(4)$  Output  
(A(4))  
Output (t) =  $H(4) + Inpu1(t)$   
Output (s) =  $H(s) \cdot Input(s)$
So...linear time-invariant systems are "simple" (not complex) for two reasons:

(1) They have the property that the impulse response fully characterizes their behavior. All responses to more complex inputs are just a convolution of the impulse response (the transfer function) with the input function.

(2) Higher order systems can always be broken down into a serial process of linked first order systems....

## Next time...a full analysis of **n = 2 linear systems**...and graphical tools

_	n = 1	n = 2 or 3	n >> 1	continuum
Linear	exponential growth and decay single step conformational change fluorescence emission pseudo first order kinetics	second order reaction kinetics linear harmonic oscillators simple feedback control sequences of conformational change	electrical circuits molecular dynamics systems of coupled harmonic oscillators equilibrium thermodynamics diffraction, Fourier transforms	Diffusion Wave propagation quantum mechanics viscoelastic systems
Nonlinear	fixed points bifurcations, multi stability irreversible hysteresis overdamped oscillators	anharmomic oscillators relaxation oscillations predator-prey models van der Pol systems Chaotic systems	systems of non- linear oscillators non-equilibrium thermodynamics protein structure/ function neural networks the cell ecosystems	Nonlinear wave propagation Reaction-diffusion in dissipative systems Turbulent/chaotic flows

Next, the obvious mathematical model....



How do we solve this differential equation? Well here is one way....

$$\frac{dB}{dt} = -k_2B + k_1A, \quad \text{where} \quad A = A_0e^{-k_1t}$$

$$B(0) = 0$$

We make a proposal....

The solution is going to be a sum of the homogeneous solution and the particular solution to the specific mport. B(+)= Bp(+) + Bh(+) Lythe particular solution.

The idea is to think of this system as having two parts to its solution....one that looks like its "natural" response (the <u>homogeneous solution</u>) and one that looks like the input into it (the <u>particular solution</u>). Let's look at the particular solution first...

The solution is going to be a sum of the homogenesies solution and  
the particular solution to the specific mpst.  
B(+)= Bp(+) + Bh(+)  
the homogene solution 
$$\frac{dB}{dE} = -k_EB$$

$$B(4) := \frac{B_{p}(4) + \frac{B_{h}(4)}{4t}}{4t}$$
 the homegon volution  $\frac{dB_{p}}{dt} := \frac{b_{q}B_{p}}{b_{q}B_{p}}$   
Now,  $B_{p}(4)$  is going to look like the import. So ... 10  
 $B_{p}(4) := C A_{0} e^{-k_{1}t}$  All we need to Biguer out is whethe C.  
 $\frac{dB_{p}}{dt} + k_{p}B_{p} = k_{1}A_{p}$   
 $= k_{1}A_{p}e^{-k_{1}t}$   
 $-k_{1}CA_{0}e^{-k_{1}t} = k_{1}A_{0}e^{-k_{1}t}$   
 $-k_{1}C + k_{2}C = k_{1}$   
 $C = \frac{k_{1}}{k_{2}-k_{1}}$   
The particular solution

And now for the "homogeneous solution"...

$$B(t) = \frac{B_{p}(t) + B_{h}(t)}{L}$$

$$= \frac{A_{0}[k_{1}]}{k_{2}-k_{1}} = \frac{B_{p}(t) + B_{h}(t)}{k_{2}-k_{1}}$$

$$= \frac{A_{0}[k_{1}]}{k_{2}-k_{1}} = \frac{k_{1}t}{t} + D e^{-k_{2}t}$$

$$= \frac{A_{0}[k_{1}]}{k_{2}-k_{1}} = \frac{k_{1}t}{t} + D$$

$$= \frac{A_{0}[k_{1}]}{k_{2}-k_{1}} + D$$

So, putting it all together...

$$B(4) = \frac{B_{p}(4) + B_{h}(4)}{L}$$
the homegee volution  $\frac{dB}{dE} = -k_{E}B$ 

$$\frac{B(4) = B_{p}(4) + B_{h}(4)$$

$$= \frac{A_{0}[k_{1}]}{k_{2}-k_{1}} = \frac{k_{1}t}{L} + D = \frac{k_{2}t}{L}$$

$$\frac{B(4) = B_{p}(4) + B_{h}(4)}{L}$$

So, putting it all together...

